



# Structure-Preserving Finite Element Methods for Fluids

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## Aim

*What is the role of conserved quantities in fluid dynamics simulations?*

## Conservation in Navier–Stokes

Let  $\Omega \subset \mathbb{R}^n$  be a domain. Incompressible Navier–Stokes:

$$\begin{aligned}\partial_t u + (u \cdot \nabla)u &= \text{Re}^{-1} \Delta u - \nabla p + f, \\ \nabla \cdot u &= 0\end{aligned}$$

with periodic or zero boundary conditions. Assuming no external force and taking “infinite Reynolds number” (zero viscosity), we get the Euler momentum equation

$$\partial_t u + (u \cdot \nabla)u = -\nabla p.$$

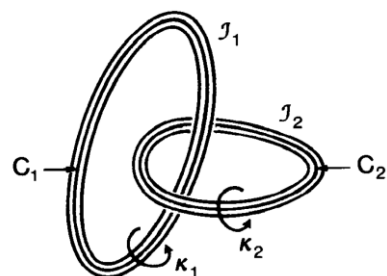
Define energy and helicity:

$$\begin{aligned}E &:= \frac{1}{2} \int_{\Omega} u \cdot u \, dx, \\ \mathcal{H} &:= \int_{\Omega} u \cdot (\nabla \times u) \, dx.\end{aligned}$$

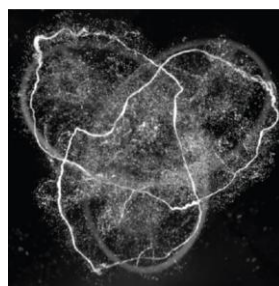
If  $u$  is a solution to incompressible Euler equations, these are **conserved in time**, i.e.  $dE/dt = 0 = d\mathcal{H}/dt$ .

What is helicity? One interpretation: quantifies “knottedness” of vortex lines: below left,  $\mathcal{H} = \pm 2n\kappa_1\kappa_2$  where  $n$  is the linking number of the vortex loops.

Aside: knotted vortex rings have been created experimentally!



Source: (1), Figure 1



Source: (2), Figure 1

## The Finite Element Method

Take Poisson problem as an example: given  $f$ , find smooth  $u$  such that

$$\begin{aligned}-\Delta u &= f \quad \text{in } \Omega, \\ u &= 0 \quad \text{on } \partial\Omega.\end{aligned}$$

Our goal: find some  $u_h$  such that  $u_h \rightarrow u$  as  $h \rightarrow 0$ .

First find variational formulation: if  $u$  solves Poisson problem then for all smooth  $v$  with value zero on  $\partial\Omega$

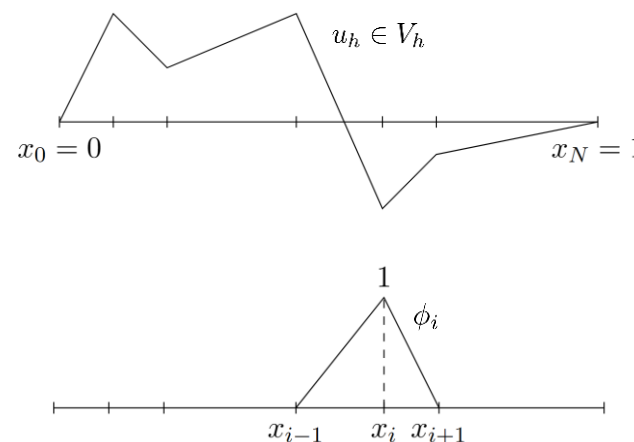
$$-\int_{\Omega} (\Delta u)v \, dx = \int_{\Omega} f v \, dx.$$

Integrate by parts to turn LHS into  $\int_{\Omega} \nabla u \cdot \nabla v \, dx$ . Now both sides are linear in  $u$  and  $v$ , so we can abstract to: find  $u \in V$  such that

$$a(u, v) = (f, v) \quad \forall v \in V.$$

For our case,  $V = H_0^1(\Omega)$ . An issue: this is infinite-dimensional, so it is difficult to solve variational problem. Solution (**Galerkin method**): replace  $V$  with a finite-dimensional  $V_h \subset V$ .

The **finite element method**: partition  $\Omega$  into simplices (intervals, triangles, tetrahedra, ...). Take  $V_h$  to consist of functions that are polynomials of some fixed degree, with “hat” basis functions:



Since  $u_h = \sum_{i=1}^d c_i \phi_i$ , Galerkin problem becomes: find  $\mathbf{c} = (c_1, \dots, c_d)$  such that

$$\sum_{i=1}^d a(\phi_i, \phi_j) c_i = (f, \phi_j) \quad \text{for all } j = 1, \dots, d,$$

i.e. solve  $A\mathbf{c} = \mathbf{F}$ .

## Preserving structures in FEM

Key point: for finite element approximation of Euler equations,  $E_h$  and  $\mathcal{H}_h$  are not necessarily conserved. However, Rebholz 2007 (3) designed numerical schemes that conserved both (replace nonlinear term with  $\omega \times u$ , modify pressure, trapezoidal timestep, solve for 3 fields).

Some further areas of research:

- Computational tradeoff? I.e. when is helicity conservation “worth it”?
- How does helicity affect dynamics (helicity cascade)?
- General manifolds (FEEC)
- Domains with nontrivial (co)homology
- Turbulence models
- Magnetohydrodynamics

## References & acknowledgements

- (1) Moffat & Tsinober, Helicity in Laminar and Turbulent Flow, *Ann. Rev. Fluid Mech.* 1992
- (2) Kleckner & Irvine, Creation and Dynamics of Knotted Vortices, *Nature Physics*, 2013
- (3) Rebholz, An Energy- and Helicity-Conserving Finite Element Scheme for the Navier–Stokes Equations, *SIAM J. Numer. Anal.* 2007

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